

Half-maximal gauged supergravities from 10d heterotic DFT

Based on: F. Hassler, Y. Sakatani, LS, *Generalized Dualities for Heterotic and Type I Strings*, 2312.16283 [hep-th], submitted to JHEP.

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07 June 2024



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Generalised dualities

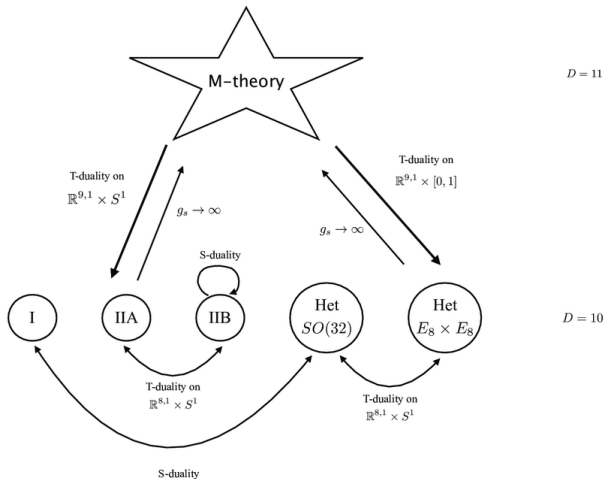


Figure – A. Fontanella, "Black Horizons and Integrability in String Theory", arXiv :1810.05434 [hep-th].

Supergravities

Superstring theories $\xrightarrow{\text{low energy limit}}$ SUGRAs.

Maximal¹ (32 supercharges)

M-theory

Type IIA

Type IIB

Half-maximal (16 supercharges)

Heterotic $E_8 \times E_8$

Heterotic $SO(32)$

Type I.

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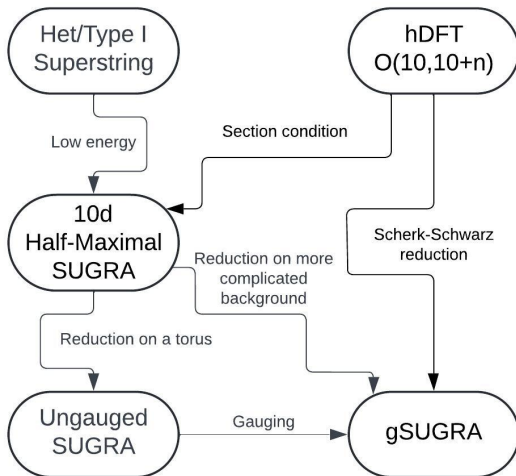
10-dim SUGRA $\xrightleftharpoons[\text{uplifts}]{\text{consistent truncations}}$ (10-d)-dim SUGRA.

Are there constraints on the procedure ?

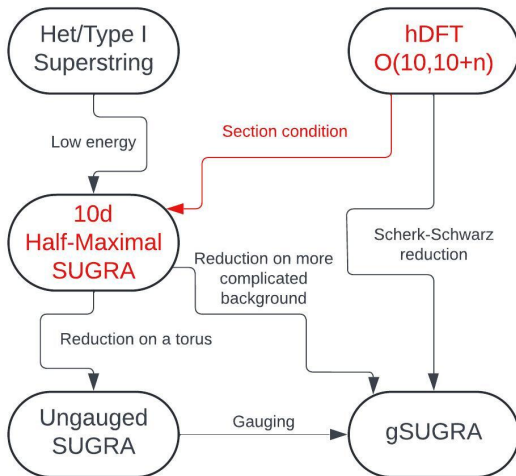
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Dimensional reductions



From hDFT to SUGRA



Heterotic Double Field Theory (hDFT)

10-dimensional hDFT : $O(10, 10 + n)$ -covariant field theory defined by² :

$$\mathcal{L}_{\text{DFT}} = e^{-2d} \left(-\frac{1}{12} H^{\text{AD}} H^{\text{BE}} H_{\text{CF}} \hat{F}_{\text{AB}}{}^{\text{C}} \hat{F}_{\text{DE}}{}^{\text{F}} - \frac{1}{4} H^{\text{AB}} \hat{F}_{\text{AD}}{}^{\text{C}} \hat{F}_{\text{BC}}{}^{\text{D}} + \right. \\ \left. + 2 H^{\text{AB}} D_{\text{A}} F_{\text{B}} + \frac{1}{6} \hat{F}_{\text{AC}}{}^{\text{D}} \hat{F}_{\text{D}}{}^{\text{A}}{}^{\text{C}} - 2 D_{\text{A}} F^{\text{A}} + F_{\text{A}} F^{\text{A}} \right), \quad (1)$$

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$\Sigma_{\text{IJ}}{}^{\text{K}}$ constant torsion term,

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From 10-dim hDFT to 10-dim SUGRA (1)

To be well defined we have to impose the constraint

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$$\Sigma_{\mathbb{I}\mathbb{J}}^{\mathbb{K}} = \begin{cases} f_I \mathcal{J}^{\mathcal{K}} & \text{structure constants of } \mathrm{Lie}(\mathcal{G}), \\ 0 & \text{otherwise.} \end{cases}$$



From 10-dim hDFT to 10-dim SUGRA (2)

Standard parametrization

$$\begin{aligned} E_{\mathbb{A}}^{\mathbb{I}} &= E_{\mathbb{A}}^{\mathbb{I}}(e, A, B_2), \\ e^{-2d} &:= e^{-2\Phi} \sqrt{-g}, \end{aligned} \tag{6}$$

we find the 10-dim half-maximal SUGRA action

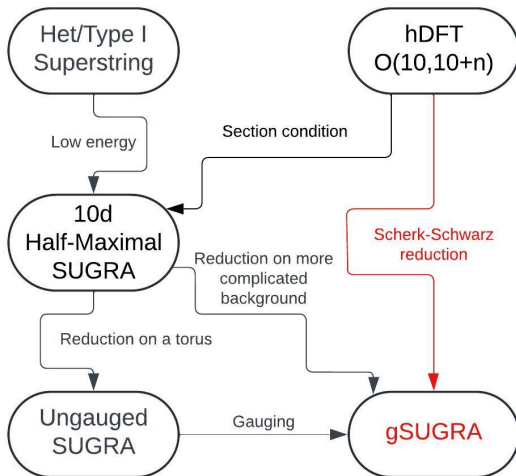
$$S = \int d^{10}x \sqrt{-g} e^{-2\Phi} \left(R + 4 \partial_m \Phi \partial^m \Phi - \frac{1}{12} \hat{H}_{mnp} \hat{H}^{mnp} - \frac{1}{4} \kappa_{I\mathcal{J}} F_{mn}{}^I F^{mn\mathcal{J}} \right), \tag{7}$$

with field strengths

$$\begin{aligned} \hat{H}_3 &:= dB_2 - \frac{1}{2} \kappa_{I\mathcal{J}} A^I \wedge dA^{\mathcal{J}} - \frac{1}{3!} f_{I\mathcal{J}\mathcal{K}} A^I \wedge A^{\mathcal{J}} \wedge A^{\mathcal{K}}, \\ F_2^I &:= dA^I + \frac{1}{2} f_{\mathcal{J}\mathcal{K}}{}^I A^{\mathcal{J}} \wedge A^{\mathcal{K}}. \end{aligned} \tag{8}$$



Half-maximal gSUGRAs



Half-maximal gSUGRAs in $10 - d$ dimensions, $1 < d \leq 6$, are classified by the embedding of their gauge group G into the Lie group³

$$G_D = \tilde{G}_D(d) \times O(d, d+n) \quad \text{with} \quad \tilde{G}(d) = \begin{cases} \mathbb{R}^+ & 0 < d \leq 5 \\ \text{SL}(2) & d = 6 \end{cases}. \quad (9)$$



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$$X_{\hat{A}\hat{B}}^{\hat{C}} \xrightarrow{\text{IRREP decomp.}} \begin{cases} \{F_{ABC}, \xi_A\} & d \leq 4, \\ \{F_{ABC}, \xi_A, \vartheta_A, \xi_{AB}, \vartheta_*\} & d = 5, \\ \{F_{\alpha ABC}, \xi_{\alpha A}, \vartheta_{\alpha A}\} & d = 6. \end{cases} \quad (10)$$

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Geometric gaugings

From G_D frames

$$\mathcal{L}_{E_{\hat{A}}} E_{\hat{B}}^{\hat{I}} = -X_{\hat{A}\hat{B}}^{\hat{C}} E_{\hat{C}}^{\hat{I}}, \quad (11)$$

we want to obtain $O(d, d+n)$ frames.

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$$\mathcal{L}_{E_A} E_B^I = -X_{AB}^C E_C^I. \quad (12)$$

$d = 5$: we have to supplement to the section condition ⁴

$$\partial_* = 0 \rightarrow \xi_{AB} = \vartheta_* = 0. \quad (13)$$

$d = 6$: we have to supplement to the section condition

$$\partial_{-I} = 0 \rightarrow F_{-ABC} = 0. \quad (14)$$

4. Y. Sakatani, *Half-maximal extended Drinfel'd algebras*, PTEP 2022 (2022) 1, 013B14.  9/12

Truncating of the 10d action

Then Scherk-Schwarz reduction to truncate.

$$F_A = e^\Delta (2 D_A \phi - \xi_A - \partial_I E_A^I), \quad F_{ABC} = X_{[ABC]}. \quad (15)$$

Computing the action it splits into

$$S = \int_{ext} d^{10-d} x \mathcal{L}_{ext} \int_{int} d^d x v, \quad (16)$$

v : scalar density \rightarrow left-invariant integration measure if the trombone gauging (\mathfrak{g} 's) vanish.

We can, then, truncate the internal part.



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- 3 we found the frames corresponding to the dimensional-reduced gSUGRAs ;
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- 5 (In the paper : generalised dualities coming from G/H construction, explicit examples, constructing frames from random gaugings and an explicit example of a generalised duality. Analysed obstructions coming from the trombone gauging.)



Thanks for your attention !

Bibliography :

- 1 F. Hassler, Y. Sakatani, LS, *Generalized Dualities for Heterotic and Type I Strings*, 2312.16283 [hep-th], submitted to JHEP.
- 2 F. Hassler, Y. Sakatani, *All maximal gauged supergravities with uplift*, PTEP 2023 (2023) 8, 083B07.
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- 4 O. Hohm, H. Samtleben, *Gauge theory of Kaluza-Klein and winding modes*, Phys.Rev.D 88 (2013), 085005.
- 5 J. Schon, M. Weidner, *Gauged $N=4$ supergravities*, JHEP 05 (2006).
- 6 E. Malek, *Half-Maximal Supersymmetry from Exceptional Field Theory*, Fortsch.Phys. 65 (2017) 10-11, 1700061.
- 7 Y. Sakatani, *Half-maximal extended Drinfel'd algebras*, PTEP 2022 (2022) 1, 013B14.



Preparing the lower dimensional reduction

We want to reduce on an internal d -dimensional space. We perform the following splitting of indices⁵

$$x^{\mathbb{I}} = \left(x^\mu \quad x^I \quad x_\mu \right), \quad \text{with} \quad x^I := \left(x^i \quad x^{\mathcal{I}} \quad x_i \right). \quad (17)$$

The decomposition reads

$$O(10, 10 + n) \rightarrow GL(10 - d) \times O(d, d + n) \quad (18)$$

$$\rightarrow GL(10 - d) \times GL(d) \times \mathcal{G}. \quad (19)$$

$$\mathbb{E}_A^{\mathbb{I}} = \begin{pmatrix} e_a^\mu & -e_a^\nu A_\nu{}^I & -e_a^\nu (B_{\nu\mu} + \frac{1}{2} A_\nu{}^K A_{\mu K}) \\ 0 & \mathcal{V}_A{}^I & \mathcal{V}_A{}^K A_{\mu K} \\ 0 & 0 & e_\mu^a \end{pmatrix}, \quad (20)$$

$$e^{-2d} = e^{-2\phi} \sqrt{-\det(g_{\mu\nu})}.$$

5. O. Hohm, H. Samtleben, *Gauge theory of Kaluza-Klein and winding modes*, Phys.Rev.D 88 (2013), 085005.



Generalized dualities

We constructed the frames

$$E_{\hat{A}}^{\hat{I}} = M_{\hat{A}}^{\hat{B}} V_{\hat{B}}^{\hat{J}} N_{\hat{J}}^{\hat{I}} \quad (21)$$

where

$$(M^{-1})_{\hat{A}}^{\hat{C}} dM_{\hat{C}}^{\hat{B}} = -v^{\hat{C}} X_{\hat{C}\hat{A}}^{\hat{B}}, \quad (22)$$

$$N_{\hat{I}}^{\hat{J}} := \left[\exp\left(-\frac{1}{2!} \mathfrak{b}_{ij} R^{ij}\right) \exp\left(-\alpha_k^I R_I^k\right) \right]_{\hat{I}}^{\hat{J}}. \quad (23)$$

From the embedding tensor we extract the $(10 - d)$ -dimensional gSUGRA gauge group G . This is the Lie group corresponding to a Lie subalgebra of the Leibniz algebra of gaugings $((T_{\hat{A}})_{\hat{B}}^{\hat{C}} = -X_{\hat{A}\hat{B}}^{\hat{C}})$

$$T_{\hat{A}} \circ T_{\hat{B}} = X_{\hat{A}\hat{B}}^{\hat{C}} T_{\hat{C}}. \quad (24)$$

We choose a subgroup H of G (usually maximally isotropic), and the frames live in G/H . Choosing different H 's result in the same physics : dualities.

Parametrization of the frame

$$\mathbb{E}_A^I := \begin{pmatrix} e^{\hat{a}_n} & 0 & 0 \\ 0 & v_{\mathcal{A}}^{\mathcal{N}} & 0 \\ 0 & 0 & e^{\hat{a}_n} \end{pmatrix} \begin{pmatrix} \delta_n^{\mathcal{P}} & -A_n^{\mathcal{P}} & -\frac{1}{2} A_n^{\mathcal{Q}} A_{\mathcal{P}\mathcal{Q}} \\ 0 & \delta_{\mathcal{N}}^{\mathcal{P}} & A_{\mathcal{P}\mathcal{N}} \\ 0 & 0 & \delta_{\mathcal{P}}^{\mathcal{N}} \end{pmatrix} \begin{pmatrix} \delta_p^m & 0 & -B_{\mathcal{P}m} \\ 0 & \delta_{\mathcal{P}}^{\mathcal{M}} & 0 \\ 0 & 0 & \delta_m^{\mathcal{P}} \end{pmatrix},$$
$$e^{-2d} := e^{-2\Phi} \sqrt{-g},$$

(25)

