

Holographic Approach to Scale vs. Conformal Invariance

Work in Progress with Lavish, Pawel Rykuta

Motivation: Conformal Group (and its Subgroups)

Algebra:

$$\left. \begin{aligned} [P_m, M_{ns}] &= 2 \eta_{m[n} P_{s]} \\ [M_{mn}, M_{rs}] &= \eta_{mr} M_{sn} - \eta_{nr} M_{sm} - \eta_{ms} M_{rn} + \eta_{ns} M_{rm} \\ [D, P_m] &= P_m \\ [K_m, M_{ns}] &= 2 \eta_{m[n} K_{s]} \\ [P_m, K_n] &= 2 (\eta_{mn} D + M_{mn}) \\ [D, K_m] &= -K_m \\ \text{Other Brackets} &= 0 \end{aligned} \right\} \begin{array}{l} \text{Poincaré inv.} \\ \\ \\ \text{Scale without} \\ \text{Conformal} \\ \text{Inv.} \\ \underline{\text{SUGRA}} \end{array}$$

Conformal

Q: Are there SwCl QFTs (2)

In $d=2$, scale implies conformal inv.

In $d > 2$: ?

See: Yu Nakayama, esp. 1302.0884

Q: Can we build SwCl models in Gauge/Gravity Duality (2)

Possible Outcomes:

QFT allows SwCl	Bulk allows SwCl	
Yes	Yes	→ Holographic Model Building
No	No	→ Holographic No-Go-Theorems
Yes	No	→ Holographic Theories are special
No	Yes	→ Forbidden Landscape

Either outcome is interesting?

Q: Can we construct a "Bulk" spacetime with Isometry group
($\sim$ Algebra of Killing vectors) $\sim$ $SO(2,1)$

Killing eq: $\mathcal{L}_K g_{\mu\nu} = \nabla_\mu K_\nu + \nabla_\nu K_\mu$

Lie-bracket: $[K_1, K_2]^\mu = K_1^\nu \partial_\nu K_2^\mu - K_2^\nu \partial_\nu K_1^\mu = K_3^\mu$

Naive Ansatz: $ds^2 = e^{2A(w)} \underbrace{(-dt^2 + dx^2)}_{\text{Minkowski}} + dz^2$

Imposing Scale Inv. fixes $A(w)$



AdS Space



Conf. Inv.

New Ansatz: (Warped) Extra dimensions in the Bulk!

Bala Subramanian: 1301.6653 ; Flory: 1711.03113

Ansatz:

(can describe
SwCl or DSI)

$$ds^2 = e^{2C(w,\theta)} \underbrace{\left[e^{\frac{2w}{L}} \overbrace{(-dt^2 + dx^2)}^{\text{Minkowski}} + dw^2 \right]}_{\text{AdS}} + e^{2B(w,\theta)} (d\theta + A_{w,\theta} dw)^2$$

$\uparrow$
 $\theta \in [0, 2\pi[$



$$ds^2 = g_{\mu\nu} dx^\mu dx^\nu = S(w,\theta) \underbrace{(-dt^2 + dx^2)}_{\text{Minkowski}} + h_{ij} dy^i dy^j$$

"Fiber"

"Base"

$y^i \in \{w, \theta\}$
 $\downarrow$
Topology of cylinder



Killing Vectors: $P_m{}^\mu, M_{nr}{}^\mu$ automatic

$$D^\mu = \begin{pmatrix} -t \\ -x \\ y^u \\ z^0 \end{pmatrix} \quad \text{because of Algebra alone.}$$

Killing eq: $\nabla_\mu \nabla_\nu + \nabla_\nu \nabla_\mu = 0$ 4d Problem for $D^\mu, g_{\mu\nu}$



$$\left. \begin{aligned} \mathcal{L}_\gamma h_{ij} &= 0 \quad (\text{Killing eq. in 2d base}) \\ \mathcal{L}_\gamma S &= 2S \end{aligned} \right\} \begin{array}{l} \text{2d Problem} \\ \text{for } \gamma^i, h_{ij}, S \end{array}$$

solution

S

$$\rightarrow f_1(G) \times S$$

h_{ij}

$$\rightarrow f_2(G) \times h_{ij}$$

for any f_1, f_2 , as long as $\mathcal{L}_\gamma G = 0$

But: We want Scale without Conf. Invariance!

→ No Killing Vector $\sim K_m^\mu$ should exist

Check Weyl tensor:

$$\mathcal{L}_K C_{\alpha\beta\gamma\delta} = 0$$

Necessary, but not sufficient
cond. for Killing vector K

$$C_{\alpha\beta\gamma\delta} = 0$$

$$\mathcal{L}_K C_{\alpha\beta\gamma\delta} = 0 \text{ trivially}$$

Killing Vector K_m can
exist.

$$C_{\alpha\beta\gamma\delta} \neq 0$$

$$C_{\alpha\beta\gamma\delta} \stackrel{\text{w.p.}}{=} \phi \cdot \tilde{C}_{\alpha\beta\gamma\delta}$$

We can show:

$$\mathcal{L}_K C_{\alpha\beta\gamma\delta} \neq 0!$$

↑
Simple, no
derivatives

$$\text{Weyl} \neq 0 \Rightarrow \underline{\underline{SwC!}}$$

Model Building:

Model: $\{S, h_{ij}, \mathcal{I}^i\}$ such that D is Killing

Conditions:

① $Weyl \neq 0 \Rightarrow SwCl$

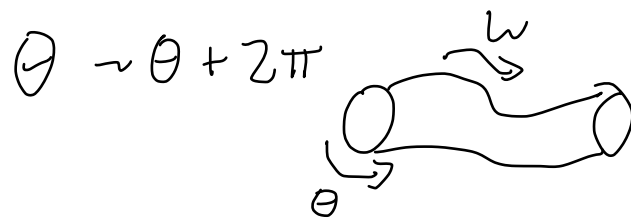
Local geometric Cond.

② NEC: $T_{\mu\nu} n^\mu n^\nu \geq 0$

Physical Cond.

③ Extra Dim. compact!

Topological Cond.



Conjecture: ① + ② + ③ mutually exclusive!

Partial Proof
Work In Progress!

NEC:

$$T_{\mu\nu} h^\mu h^\nu \geq 0 \quad \Rightarrow \quad R_{\mu\nu} h^\mu h^\nu \geq 0$$

Einstein eq.

$$g_{\mu\nu}^{\text{w.p.}} = \begin{pmatrix} \overbrace{S \eta_{ab}}^{2 \times 2 \text{ Fiber}} & 0 \\ 0 & \underbrace{h_{ij}}_{2 \times 2 \text{ Base}} \end{pmatrix} \Rightarrow R_{\mu\nu}^{\text{w.p.}} = \begin{pmatrix} X \cdot \eta_{ab} & 0 \\ 0 & r_{ij} \end{pmatrix}$$

$$h^\mu = \begin{pmatrix} \hat{n}^a \\ \bar{n}^i \end{pmatrix} \begin{matrix} \} \text{Fiber} \\ \} \text{Base} \end{matrix}$$

$$h^\mu h_\mu = 0 \Rightarrow \hat{n}^a \hat{n}^b \eta_{ab} = - \frac{\bar{n}^i \bar{n}^j h_{ij}}{S}$$

$$\Rightarrow \text{NEC: } 0 \leq R_{\mu\nu} h^\mu h^\nu = \bar{n}^i \bar{n}^j \underbrace{\left(r_{ij} - \frac{X}{S} h_{ij} \right)}_{\equiv t_{ij}}$$

$$NEC \Leftrightarrow t^i_j \text{ is a positive semi-def. } 2 \times 2 \text{ matrix}$$

$$\text{Tr} [t^i_j] = \bar{R} + \frac{2}{S} (\bar{\nabla} \sqrt{S})^2 \quad \bar{R} = R[h] \text{ in } 2d'$$

$$\text{Det} [t^i_j] = \frac{1}{4} \left[\bar{R} + \frac{2}{S} (\bar{\nabla} \sqrt{S})^2 \right]^2 + \frac{1}{S} (\bar{\nabla} \sqrt{S})^2 - \frac{2}{S} (\bar{\nabla}_i \bar{\nabla}_j \sqrt{S}) (\bar{\nabla}^i \bar{\nabla}^j \sqrt{S})$$

Naive Idea: just make $\bar{R} > 0$ everywhere (and $\bar{R}$ big enough) to guarantee NEC
 $(\rightarrow \text{det} > 0, \text{tr} > 0)$

But: We can compactify cylinder to torus b.c. of Killing vector ∂_θ :



Gauss-Bonnet: $\int \bar{R} dV \sim 2\pi \chi = 0 \Rightarrow \bar{R} > 0$ everywhere is not possible!

Outlook / To Do :

— Finalise Proof

— What if extre dim is not compact?
E.g. defect / janus like models?

— How much changes as dimension
of base and fiber changes?

Paper
Soon

Thank you very much!